

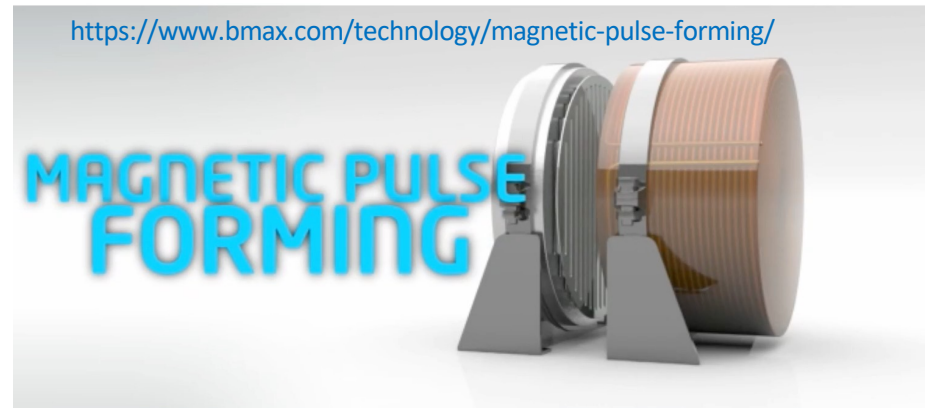
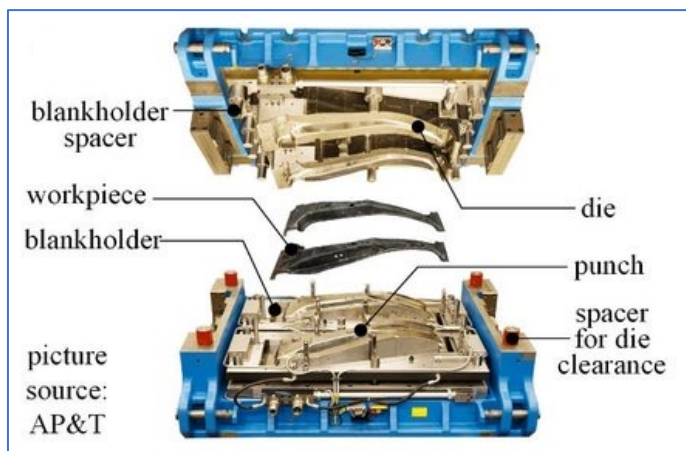
Force and stress tensor

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Intro

- 재료에 가해질 수 있는 힘은
 - 1) 재료 표면에 가해진 '접촉'에 의해 전달되는 힘 (surface force)
 - 2) 접촉하지 않고서 전달되는 힘 (body force)
 - 중력, 전자기력 등
- 우리가 '정역학'에서 다루는 재료의 '운동'상태 혹은 '변형'은 모두 평형 상태인 '힘'들에 의해서 나타난다. (평형상태가 아닌 힘에 의해 발생하는 현상은 다루지 않음)
- 일반적인 소성가공 공정에서 재료는 surface force에 의해서만 힘을 전달 받고 변형하며, body force에 의한 소성 가공법은 매우 특별한 경우에 한정되어 있다.

Source:AP&T, (<https://www.aptgrou.com>)



Magnetic pulse forming
<https://youtu.be/rBXXBIP9qIE>

Stamping process for mass production



TESLA Stamping line: <https://youtu.be/gkjin9bogLSM>



Force

Force is a vector quantity, so that

$$\mathbf{f} = f_1 \mathbf{e}_1 + f_2 \mathbf{e}_2 + f_3 \mathbf{e}_3$$

Force equilibrium condition:

$$\sum_{\text{source of acting force}} \mathbf{f}^{\text{each source}} = 0$$

$$\sum_{\text{source}} f_i = 0 \quad (\text{free index } i) \rightarrow \sum_{\text{source}} f_1 = 0, \sum_{\text{source}} f_2 = 0, \sum_{\text{source}} f_3 = 0$$

Force and moment

$\mathbf{m} = \mathbf{f} \times \mathbf{x}$ with \mathbf{x} being moment arm

Moment of a force is a measure of its tendency to cause a body to 'rotation' about a specific point or axis.

$$\begin{aligned}\mathbf{m}_k &= f_i x_j \mathbf{e}_i \times \mathbf{e}_j = f_i x_j \epsilon_{ijk} \mathbf{e}_k \\ &= \mathbf{e}_k \sum_i \sum_j f_i x_j \epsilon_{ijk} \\ &= \mathbf{e}_k \{f_1 x_2 \epsilon_{12k} + f_1 x_3 \epsilon_{13k} + f_2 x_1 \epsilon_{21k} + f_2 x_3 \epsilon_{23k} + f_3 x_1 \epsilon_{31k} + f_3 x_2 \epsilon_{32k}\}\end{aligned}$$

$$\begin{aligned}\mathbf{m}_1 &= \mathbf{e}_1 \{f_1 x_2 \epsilon_{121} + f_1 x_3 \epsilon_{131} + f_2 x_1 \epsilon_{211} + f_2 x_3 \epsilon_{231} + f_3 x_1 \epsilon_{311} + f_3 x_2 \epsilon_{321}\} \\ &= \mathbf{e}_1 \{f_2 x_3 - f_3 x_2\}\end{aligned}$$

$$\begin{aligned}\mathbf{m}_2 &= \mathbf{e}_2 \{f_1 x_2 \epsilon_{122} + f_1 x_3 \epsilon_{132} + f_2 x_1 \epsilon_{212} + f_2 x_3 \epsilon_{232} + f_3 x_1 \epsilon_{312} + f_3 x_2 \epsilon_{322}\} \\ &= \mathbf{e}_2 \{-f_1 x_3 + f_3 x_1\}\end{aligned}$$

$$\mathbf{m}_3 = \mathbf{e}_3 \{f_1 x_2 - f_2 x_1\}$$

$m_i = 0$ (moment equilibrium; i is a free index)

응력 텐서 (stress tensor)

- 재료에 전달되는 '힘'을 '세기 물리량'으로 정량화하여 나타내는데 '응력'이라는 물리량을 사용한다.
- 응력은 '텐서' 물리량이다.
- 응력 텐서의 정의를 자세히 알아보자.
 - 힘 평형 (force equilibrium)
 - 응력 벡터 (stress vector)

힘 벡터와 응력 벡터

힘은 크기 물리량; 응력은 세기 물리량

풍선에게 동일한 힘을 손바닥으로 전달할 때와 바늘끝으로 전달할 때 풍선에게 나타날 효과가 매우 다른 것을 예상할 수 있다. 힘이 작용하는 면적의 차이에 의해서 나타나는 효과라 볼 수 있다.



• 응력 벡터(혹은 traction이라 일컫는다)란? $\mathbf{t} = \lim_{\Delta S \rightarrow 0} \frac{\mathbf{f}}{\Delta S}$,

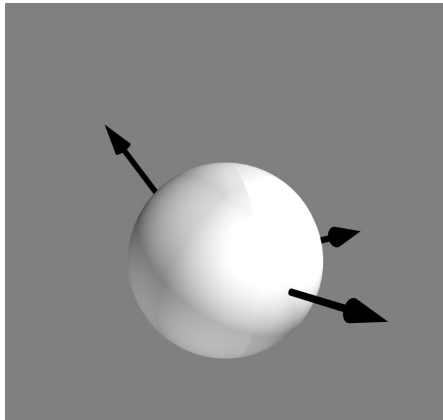
• In SI unit system, unit of $\frac{f}{\Delta S} = \frac{\left[kg \frac{m}{s^2} \right]}{[m^2]} = \frac{N}{m^2} = Pa$

Caution1 : 변형률은 unit이 없지만, 응력은 단위가 있다.

Caution2: 금속에 작용하는 응력은 주로 MPa가 쓰이며, 여기서 M은 Mega이며 뜻은 10^6 이다.

The graphical illustration to understand the definition of stress vector (traction vector)

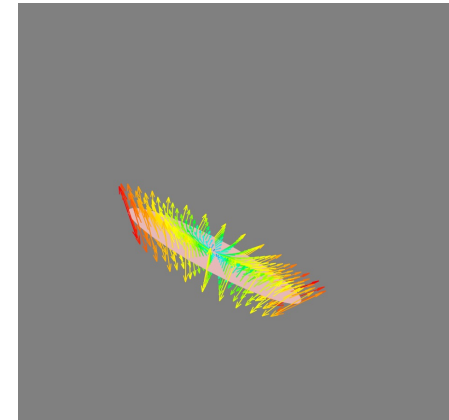
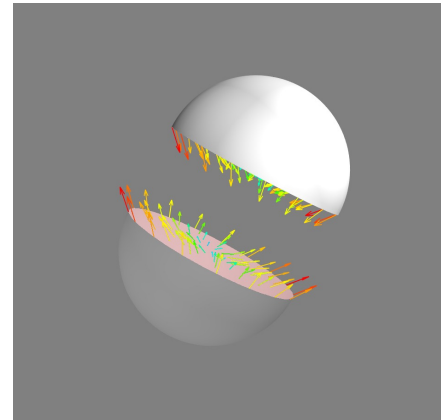
외부의 surface force



재료 내부에 force field
 $f(x, y, z)$



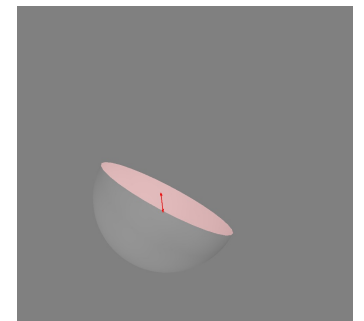
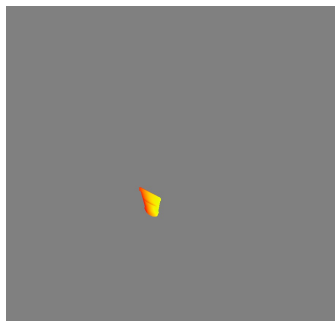
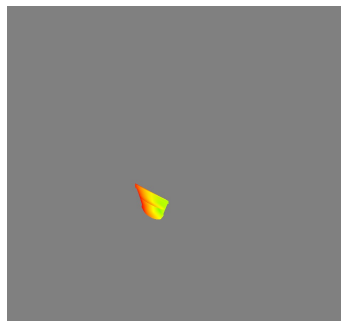
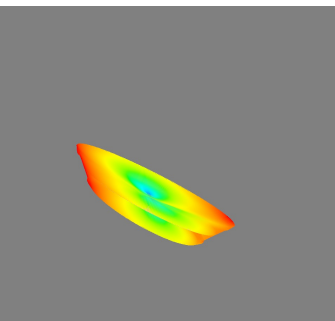
힘평형을 이론채로
force field 발달



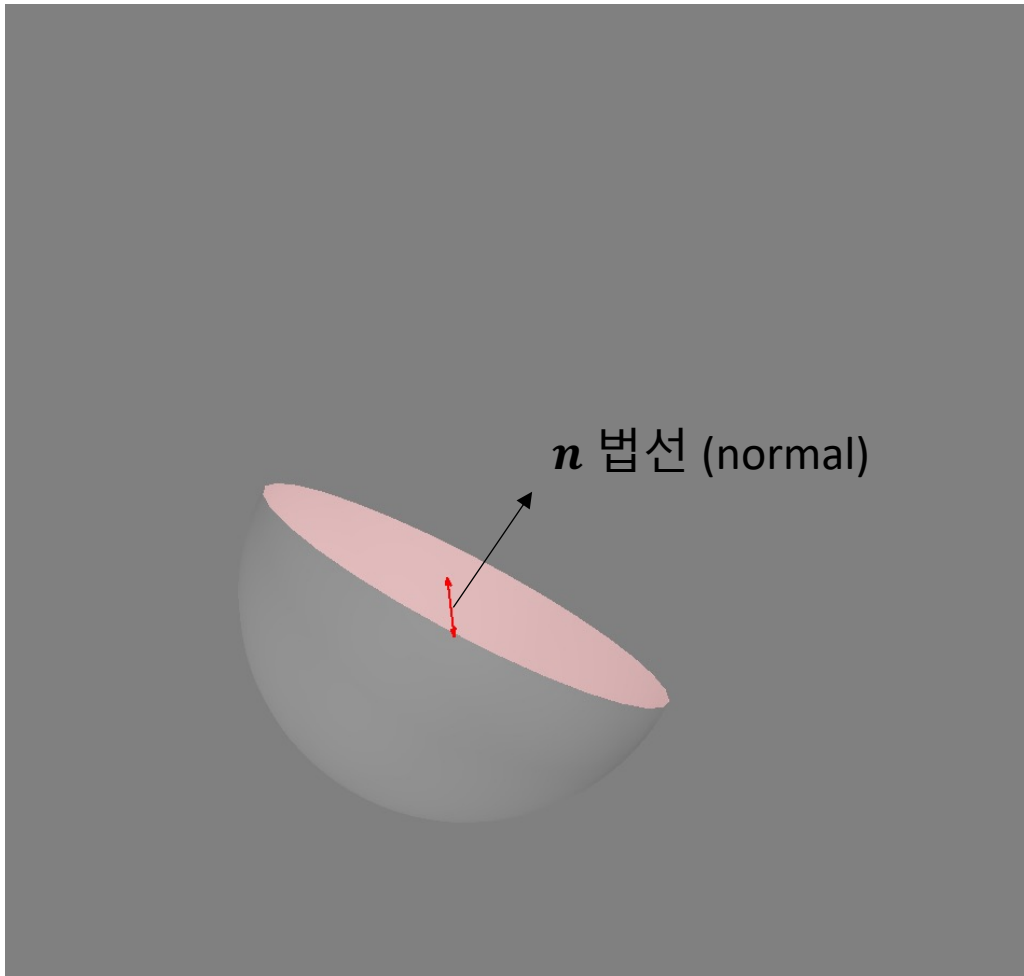
좀 더 빠르게 나타낸
힘벡터 분포

$$\mathbf{t} = \lim_{\Delta S \rightarrow 0} \frac{\mathbf{f}}{\Delta S}$$

\mathbf{t} : normal이 \mathbf{n} 인 면에
작용하는 응력 벡터



The graphical gist of it

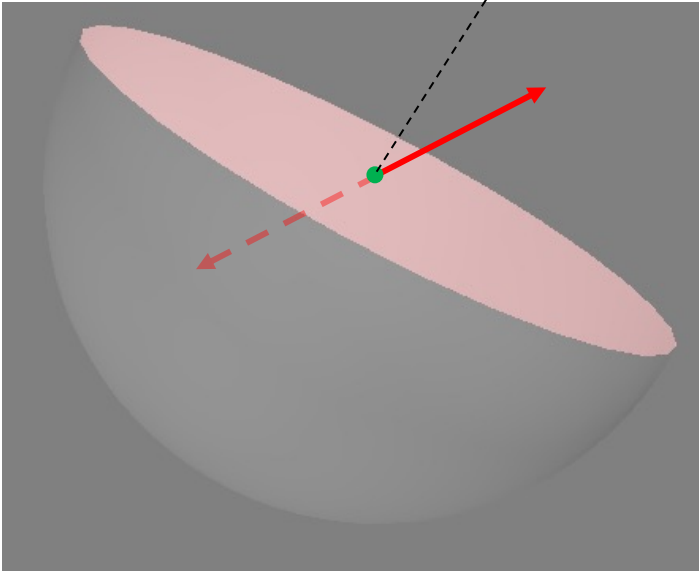


$$t = \lim_{\Delta S \rightarrow 0} \frac{f}{\Delta S}$$

$$t = \frac{df}{dS}$$

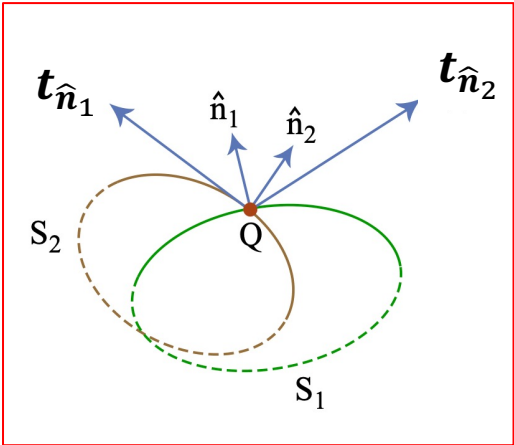
$$t_i = \frac{df_i}{dS}$$

Mathematical description (both algebraic and geometrical)

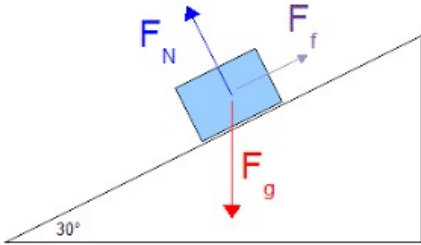


$$\boldsymbol{t} = \lim_{\Delta S \rightarrow 0} \frac{\boldsymbol{f}}{\Delta S}$$

$$\boldsymbol{t} = \frac{d\boldsymbol{f}}{dS} \qquad t_i = \frac{df_i}{dS}$$




물질내 고정된 한 점에서의 traction은 그 점을 통과하는 면의 방향에 따라 달라진다. 즉 **t**는 **n**에 따라 달라진다.




$$\boldsymbol{t} = \boldsymbol{\sigma} \cdot \boldsymbol{n}$$

$$t_i = \sigma_{ij}n_j$$

● The infinitesimal surface that approaches to zero (i.e., $\Delta S \rightarrow 0$) Stress tensor ($\boldsymbol{\sigma}$) on that location (●) linearly transforms the normal \boldsymbol{n} to the traction vector \boldsymbol{t}

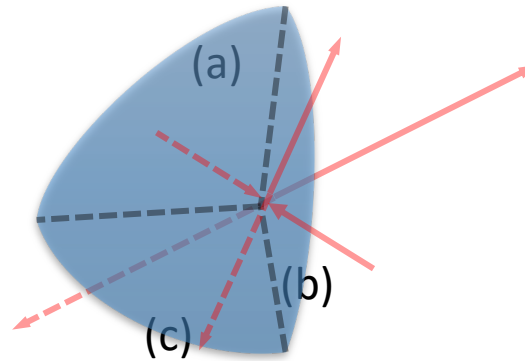
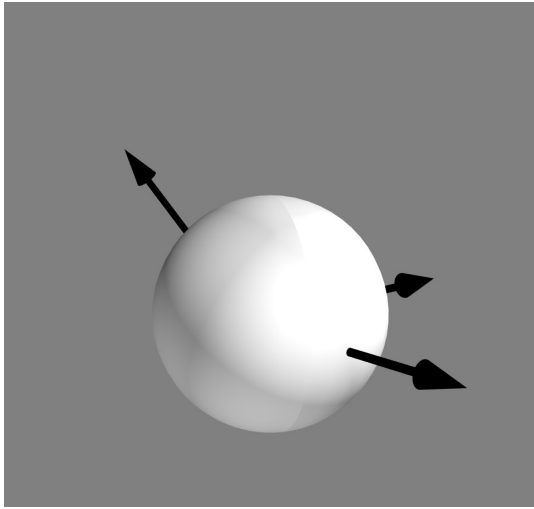
 (Coupled) Force vectors acting on that infinitesimal surface



 Plane normal; It is a unit vector such that $|\boldsymbol{n}| = 1$.

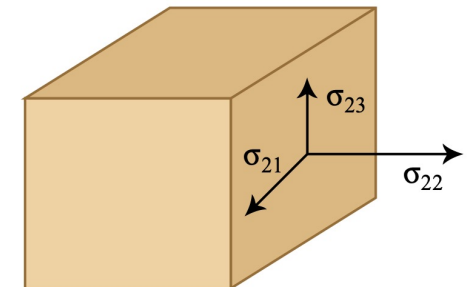
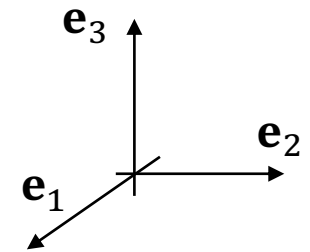
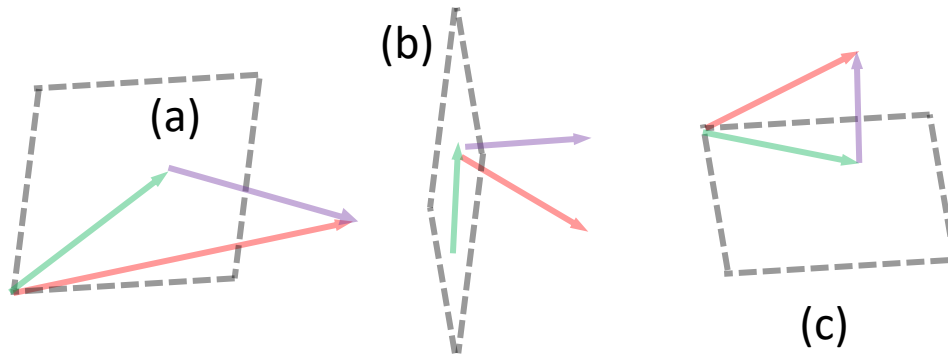
$$\begin{bmatrix} t_1 \\ t_2 \\ t_3 \end{bmatrix} = \boldsymbol{\sigma} \cdot \begin{bmatrix} n_1 \\ n_2 \\ n_3 \end{bmatrix}$$


$$\begin{bmatrix} t_1 \\ t_2 \\ t_3 \end{bmatrix} = \begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{bmatrix} \begin{bmatrix} n_1 \\ n_2 \\ n_3 \end{bmatrix}$$

Mathematical description (both algebraic and geometrical)



 Normal component
 Tangential component




 각 면에 '누워서' (tangential) 작용하는 성분을 다시 \mathbf{e}_i 에 projection 시켜 shear term은 2개, normal component는 1개가 각 면마다 있음.

Let's say by $\mathbf{t}_n = \frac{d\mathbf{f}}{dS}$

Here, we denote traction vector \mathbf{t} acting on infinitesimal surface dS with the normal \mathbf{n} as \mathbf{t}_n

$\mathbf{t}_{\mathbf{e}_i}$?

Traction vector \mathbf{t} acting on infinitesimal surface dS with the basis normal \mathbf{e}_i (note that i is a free index, it could be 1 or 2 or 3)

Three traction vectors on three planes

$$\mathbf{t}_{\mathbf{e}_1} = (\sigma_{11}, \sigma_{12}, \sigma_{13})$$

$$\mathbf{t}_{\mathbf{e}_2} = (\sigma_{21}, \sigma_{22}, \sigma_{23})$$

$$\mathbf{t}_{\mathbf{e}_3} = (\sigma_{31}, \sigma_{32}, \sigma_{33})$$

$$\mathbf{t} = \boldsymbol{\sigma} \cdot \mathbf{n}$$

$$t_i = \sigma_{ij} n_j$$

$$\mathbf{t}_{\mathbf{e}_1} = \boldsymbol{\sigma} \cdot \mathbf{e}_1 \quad ?$$

$$\begin{bmatrix} t_1 \\ t_2 \\ t_3 \end{bmatrix} = \begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{bmatrix} \begin{bmatrix} n_1 \\ n_2 \\ n_3 \end{bmatrix}$$

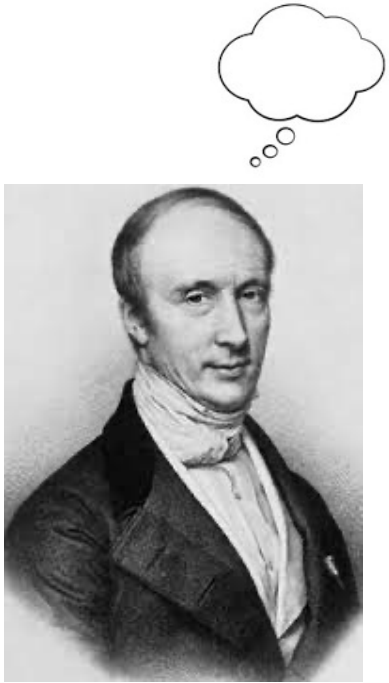
$$\begin{bmatrix} (\mathbf{t}_{\mathbf{e}_1})_1 \\ (\mathbf{t}_{\mathbf{e}_1})_2 \\ (\mathbf{t}_{\mathbf{e}_1})_3 \end{bmatrix} = \begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{bmatrix} \begin{bmatrix} (\mathbf{e}_1)_1 \\ (\mathbf{e}_1)_2 \\ (\mathbf{e}_1)_3 \end{bmatrix}$$

$$\begin{bmatrix} (\mathbf{t}_{\mathbf{e}_1})_1 \\ (\mathbf{t}_{\mathbf{e}_1})_2 \\ (\mathbf{t}_{\mathbf{e}_1})_3 \end{bmatrix} = \begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{bmatrix} \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} (\mathbf{t}_{\mathbf{e}_1})_1 \\ (\mathbf{t}_{\mathbf{e}_1})_2 \\ (\mathbf{t}_{\mathbf{e}_1})_3 \end{bmatrix} = \begin{bmatrix} \sigma_{11} \\ \sigma_{21} \\ \sigma_{31} \end{bmatrix}$$

Caution: $\sigma_{ij} = \sigma_{ji}$

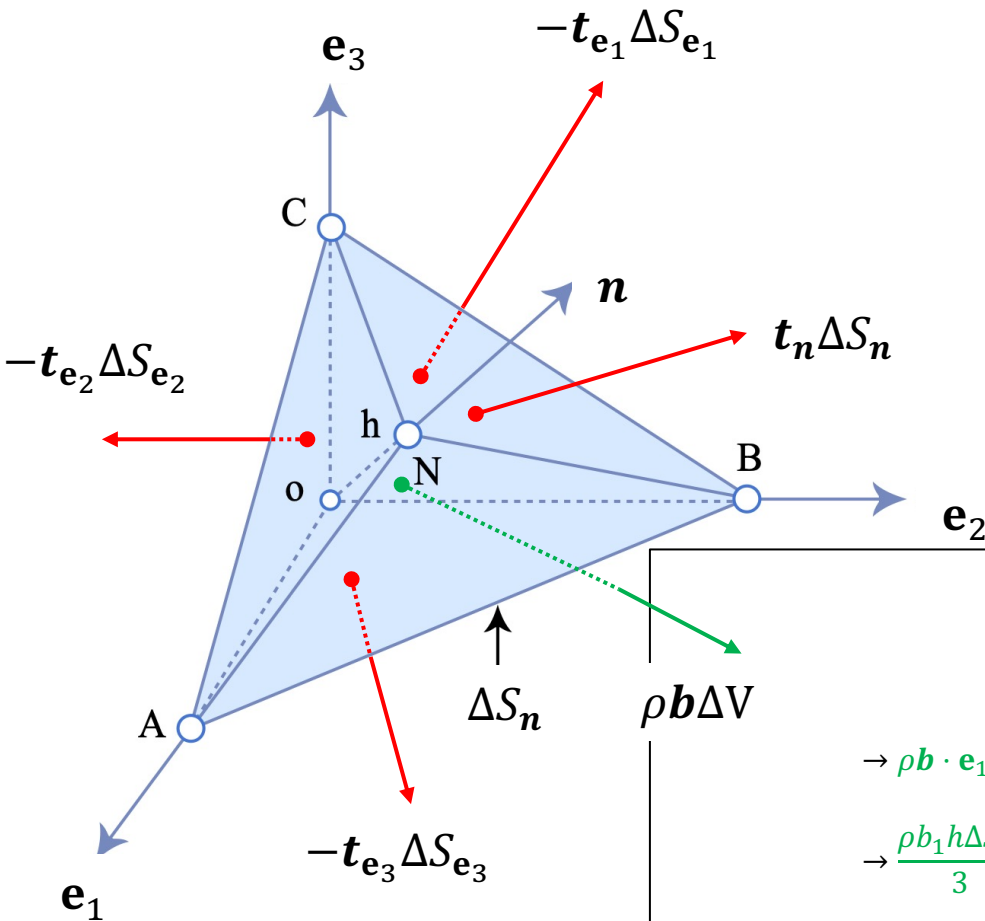
Cauchy stress tensor



Augustin-Louis Cauchy

- Cauchy's idea: Traction vectors on three independent (perpendicular) planes pertaining to a point would suffice to provide the stress state of that point.
- Components of traction vectors on each of these three planes provide us the Cauchy stress components, which is sufficient to describe any arbitrary stress state. (No more than three planes!)
- While there are a number of strain measures ...

Cauchy's Tetrahedron (사면체)



$$\mathbf{t} = \boldsymbol{\sigma} \cdot \mathbf{n}$$

$$t_i = \sigma_{ij} n_j$$

$$\mathbf{t}_{\mathbf{e}_1} = \boldsymbol{\sigma} \cdot \mathbf{e}_1$$

- \mathbf{e}_i : basis vectors (i is free)
- $\mathbf{t}_{\mathbf{e}_i}$: Traction acting on the plane $\perp \mathbf{e}_i$
- ΔS_n : Area of A,B,C 삼각형
- $\Delta S_{\mathbf{e}_i}$: surface of the plane $\perp \mathbf{e}_i$, which can be obtained by $\Delta S_{\mathbf{e}_i} = \Delta S_n \mathbf{n} \cdot \mathbf{e}_i = \Delta S_n n_i$ (Notice n_i is not in bold-face)
- b_i is the body force
- $m = \frac{1}{3} \rho h \Delta S_n$

Force balance along \mathbf{e}_1 ?

Check $F_1 = ma_1$ (net F_1)

$$\mathbf{e}_1 \cdot \Sigma \mathbf{F} = \frac{1}{3} \rho h \Delta S_n \mathbf{a} \cdot \mathbf{e}_1$$

$$\rightarrow \rho \mathbf{b} \cdot \mathbf{e}_1 \Delta V + (t_n \Delta S_n) \cdot \mathbf{e}_1 + \left(\sum_i -\mathbf{t}_{\mathbf{e}_i} \Delta S_{\mathbf{e}_i} \right) \cdot \mathbf{e}_1 = \frac{1}{3} \rho h \Delta S_n a_1$$

$$\rightarrow \frac{\rho b_1 h \Delta S_n}{3} + t_n \cdot \mathbf{e}_1 \Delta S_n - \Delta S_n \left\{ \sum_i \mathbf{t}_{\mathbf{e}_i} n_i \right\} \cdot \mathbf{e}_1 = \frac{1}{3} \rho h \Delta S_n a_1$$

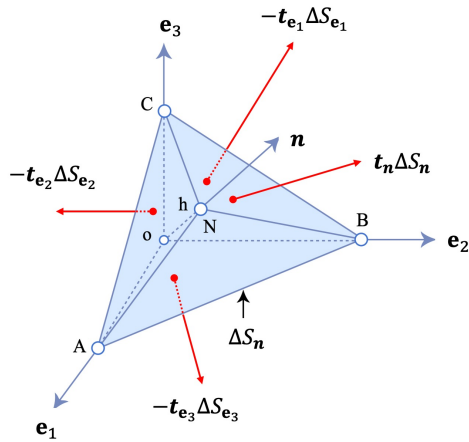
$$(LHS) \rightarrow \frac{\rho b_1 h \Delta S_n}{3} + t_n \cdot \mathbf{e}_1 \Delta S_n - \Delta S_n \left\{ \sum_i n_i \boldsymbol{\sigma} \cdot \mathbf{e}_i \right\} \cdot \mathbf{e}_1$$

$$(LHS) \rightarrow \frac{\rho b_1 h \Delta S_n}{3} + t_n \cdot \mathbf{e}_1 \Delta S_n - \Delta S_n \{ n_1 \boldsymbol{\sigma} \cdot \mathbf{e}_1 + n_2 \boldsymbol{\sigma} \cdot \mathbf{e}_2 + n_3 \boldsymbol{\sigma} \cdot \mathbf{e}_3 \} \cdot \mathbf{e}_1$$

$$(LHS) \rightarrow \frac{\rho b_1 h \Delta S_n}{3} + t_n \cdot \mathbf{e}_1 \Delta S_n - \Delta S_n \{ n_1 \boldsymbol{\sigma} \cdot \mathbf{e}_1 \cdot \mathbf{e}_1 + n_2 \boldsymbol{\sigma} \cdot \mathbf{e}_2 \cdot \mathbf{e}_1 + n_3 \boldsymbol{\sigma} \cdot \mathbf{e}_3 \cdot \mathbf{e}_1 \}$$

$$F_1 = ma_1 \rightarrow \frac{\rho b_1 h \Delta S_n}{3} + t_n \cdot \mathbf{e}_1 \Delta S_n - \Delta S_n \{ n_1 \sigma_{11} + n_2 \sigma_{21} + n_3 \sigma_{31} \} = \frac{1}{3} \rho h \Delta S_n a_1$$

Cauchy's Tetrahedron (사면체)



- \mathbf{e}_i : basis vectors
- $\mathbf{t}_{\mathbf{e}_i}$: Traction acting on the plane $\perp \mathbf{e}_i$
- ΔS_n : Area of A,B,C 삼각형
- $\Delta S_{\mathbf{e}_i}$: surface of the plane $\perp \mathbf{e}_i$, which can be obtained by $\Delta S_{\mathbf{e}_i} = \Delta S_n \mathbf{n} \cdot \mathbf{e}_i = \Delta S_n n_i$ (Notice n_i is not in bold-face)
- b_i is the body force
- $m = \frac{1}{3} \rho h \Delta S_n$

Force balance along \mathbf{e}_1 ?

Check $F_1 = ma_1$ (net F_1)

$$F_1 = ma_1 \rightarrow \frac{\rho b_1 h \Delta S_n}{3} + \mathbf{t}_n \cdot \mathbf{e}_1 \Delta S_n - \Delta S_n \{n_1 \sigma_{11} + n_2 \sigma_{21} + n_3 \sigma_{31}\} = \frac{1}{3} \rho h \Delta S_n a_1$$

With $h \rightarrow 0$

$$\rightarrow \mathbf{t}_n \cdot \mathbf{e}_1 - \{n_1 \sigma_{11} + n_2 \sigma_{21} + n_3 \sigma_{31}\} = \frac{1}{3} \rho h a_1$$

$$\mathbf{t}_n \cdot \mathbf{e}_1 = \sum_j \sigma_{j1} n_j$$

Let's notice that $\mathbf{t}_n \cdot \mathbf{e}_1 = (\mathbf{t}_n)_1$ and maybe it's better for us to shorten the notation $(\mathbf{t}_n)_1$ simply to t_1

Cauchy's Tetrahedron (사면체) (Advanced)

A case without body force

$$\int_V \nabla \phi dV = \int_S \mathbf{n} \phi dS$$

- Say, ϕ is a quantity that is transferable (like heat, momentum, solute in solution etc.), transfer-in and -out should be met if equilibrium;
- In equilibrium (such as heat equilibrium or force equilibrium), $\phi = \text{const.}$
- Therefore, (\mathbf{n}_S denotes normal of surface S)

$$\int_V \nabla \phi dV = \int_S \mathbf{n}_S \phi dS \rightarrow 0 = \int_S \mathbf{n}_S \phi dS \rightarrow 0 = \int_S \mathbf{n}_S dS$$

- For the case of this tetrahedron we have

$$\int_S \mathbf{n}_S dS = 0 \rightarrow \mathbf{n} \Delta S_n - \mathbf{e}_1 \Delta S_{e_1} - \mathbf{e}_2 \Delta S_{e_2} - \mathbf{e}_3 \Delta S_{e_3} = 0$$

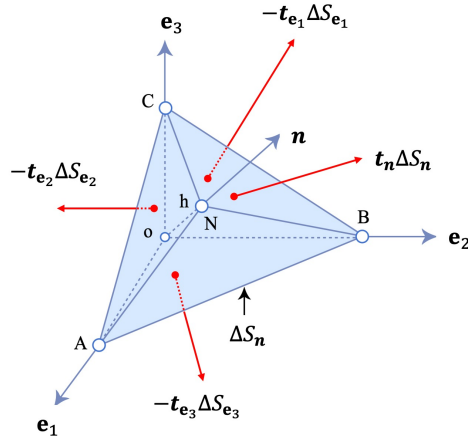
$$\rightarrow \mathbf{n} \Delta S_n = \sum_i \mathbf{e}_i \Delta S_{e_i} \rightarrow (\mathbf{n} \Delta S_n) \cdot \mathbf{e}_k = \mathbf{e}_k \cdot \sum_i \mathbf{e}_i \Delta S_{e_i}$$

$$\rightarrow n_k \Delta S_n = \sum_i \mathbf{e}_k \cdot \mathbf{e}_i \Delta S_{e_i} = \Delta S_{e_k} \rightarrow n_i \Delta S_n = \Delta S_{e_i}$$

$$\rightarrow \text{Force equilibrium: } \mathbf{t}_n \Delta S_n - \sum_i \mathbf{t}_{e_i} n_i \Delta S_n = 0 \rightarrow \mathbf{t}_n - \sum_i \mathbf{t}_{e_i} n_i = 0$$

$$\mathbf{t}_n = \sum_i \mathbf{t}_{e_i} n_i \rightarrow \mathbf{t}_n = (\mathbf{t}_{e_1} n_1 + \mathbf{t}_{e_2} n_2 + \mathbf{t}_{e_3} n_3)$$

$$\rightarrow \mathbf{t}_n = \{(\mathbf{n} \cdot \mathbf{e}_1) \mathbf{t}_{e_1} + (\mathbf{n} \cdot \mathbf{e}_2) \mathbf{t}_{e_2} + (\mathbf{n} \cdot \mathbf{e}_3) \mathbf{t}_{e_3}\}$$



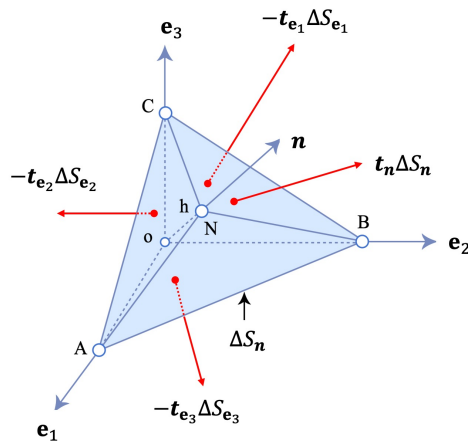
Force equilibrium:

$$\mathbf{t}_n \Delta S_n - \sum_i \mathbf{t}_{e_i} \Delta S_{e_i} = 0$$

Remember: $S_{e_i} = \Delta S_n \mathbf{n} \cdot \mathbf{e}_i = \Delta S_n n_i$

Cauchy's Tetrahedron (사면체) (Advanced)

A case without body force



$$\mathbf{t}_n = \{(\mathbf{n} \cdot \mathbf{e}_1)\mathbf{t}_{e_1} + (\mathbf{n} \cdot \mathbf{e}_2)\mathbf{t}_{e_2} + (\mathbf{n} \cdot \mathbf{e}_3)\mathbf{t}_{e_3}\}$$

$$\{t_n\}_i = \{(\mathbf{n} \cdot \mathbf{e}_1)\mathbf{t}_{e_1} + (\mathbf{n} \cdot \mathbf{e}_2)\mathbf{t}_{e_2} + (\mathbf{n} \cdot \mathbf{e}_3)\mathbf{t}_{e_3}\}_i ; \text{ note the free index } i$$

$$\{t_n\}_i = (\mathbf{n} \cdot \mathbf{e}_1)\{t_{e_1}\}_i + (\mathbf{n} \cdot \mathbf{e}_2)\{t_{e_2}\}_i + (\mathbf{n} \cdot \mathbf{e}_3)\{t_{e_3}\}_i$$

Notice that both $\{t_n\}_i$ and $\{t_{e_k}\}_i$ (k is free index) are 'scalars' – they for each index represent each component, as in below,

$$c\mathbf{a} = ca_i\mathbf{e}_i = a_i(c\mathbf{e}_i)$$

so that

$$\begin{aligned} \{t_n\}_i &= (\mathbf{n}) \cdot \mathbf{e}_1\{t_{e_1}\}_i + (\mathbf{n}) \cdot \mathbf{e}_2\{t_{e_2}\}_i + (\mathbf{n}) \cdot \mathbf{e}_3\{t_{e_3}\}_i \\ &= \mathbf{n} \cdot \{\mathbf{e}_1\mathbf{t}_{e_1} + \mathbf{e}_2\mathbf{t}_{e_2} + \mathbf{e}_3\mathbf{t}_{e_3}\}_i \end{aligned}$$

$$\mathbf{t}_n \cdot \mathbf{e}_i = \mathbf{n} \cdot (\mathbf{e}_1\mathbf{t}_{e_1} + \mathbf{e}_2\mathbf{t}_{e_2} + \mathbf{e}_3\mathbf{t}_{e_3}) \cdot \mathbf{e}_i$$

Force equilibrium:

$$t_n\Delta S_n - \sum_i t_{e_i}\Delta S_{e_i} = 0$$

Remember: $S_{e_i} = \Delta S_n \mathbf{n} \cdot \mathbf{e}_i = \Delta S_n n_i$



"I'm a college professor, Jason. You need to ask someone else if you want advice about the real world."

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Stress tensor is symmetric

$$\sigma_{ij} = \sigma_{ji}$$

Therefore, the below

$$\mathbf{t}_n \cdot \mathbf{e}_1 = \sum_j \sigma_{j1} n_j$$

could be equivalently written as:

$$(\mathbf{t}_n)_1 = \sum_j \sigma_{1j} n_j$$

And, therefore, by making use Indicinal and Einstein notations,

$$(\mathbf{t}_n)_i = \sigma_{ij} n_j$$

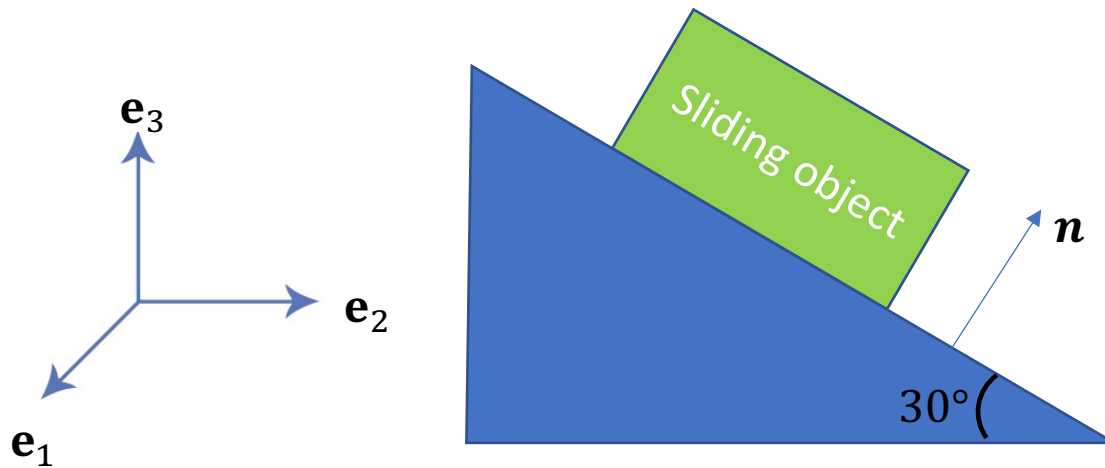
Proof of symmetric stress tensor

The same linear transformation applies to **strain** tensor

$\boldsymbol{\varepsilon} \cdot \boldsymbol{n}$ gives you the 'stretched' vector.

If you know the meaning of dot product, you'll understand $(\boldsymbol{\varepsilon} \cdot \boldsymbol{n}) \cdot \boldsymbol{n}$ gives you a fractional change in length expected along \boldsymbol{n} expected as an outcome of $\boldsymbol{\varepsilon}$.

Object sliding downhill



Points under the object is under the stress of

$$\begin{bmatrix} 0 & 0 & 0 \\ 0 & -25 & 0 \\ 0 & 0 & 40 \end{bmatrix}$$

Q: What is the traction vector acting on the plane?

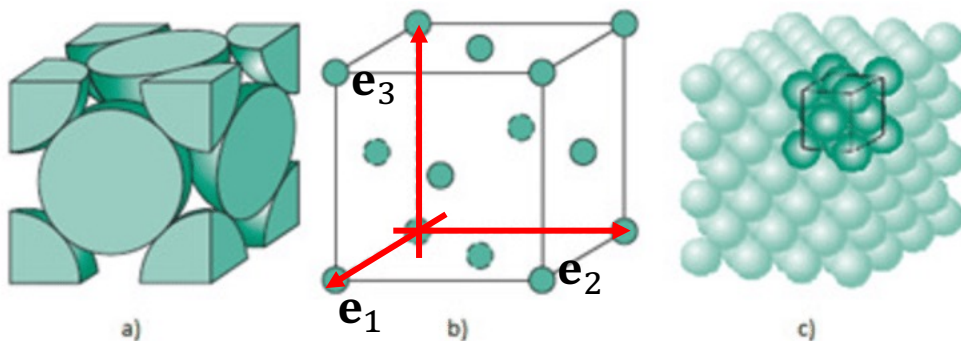
solution:

$$n = [0, \cos 60^\circ, \sin 60^\circ] = \left[0, 0.5, \frac{\sqrt{3}}{2}\right]$$

$$t = \sigma \cdot n \rightarrow \begin{bmatrix} 0 & 0 & 0 \\ 0 & -25 & 0 \\ 0 & 0 & -40 \end{bmatrix} \begin{bmatrix} 0 \\ 0.5 \\ 0.866 \end{bmatrix} = [0, -12.5, -34.6]$$

Examples.

Suppose a single crystal Ni (in FCC) is subjected to a traction field of $(100, 10, 10)$ MPa.



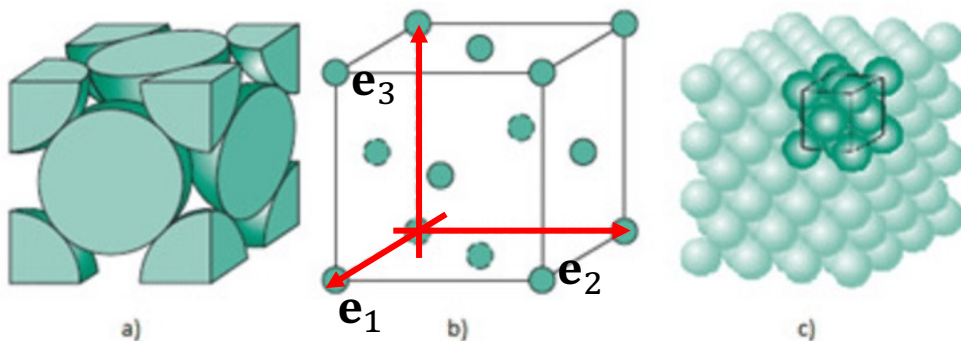
Calculate the traction vector along (111) plane normal

Solution). Since, say, $\mathbf{n} = \frac{(1,1,1)}{|(1,1,1)|} = \frac{1}{\sqrt{3}}(1,1,1)$. The component of traction vector (give as $\mathbf{t} = (100,10,10)$), its component along \mathbf{n} is obtained by projecting \mathbf{t} along \mathbf{n} such that

$$\mathbf{t} \cdot \mathbf{n} = \frac{120}{\sqrt{3}}$$

Examples.

Suppose a single crystal Ni (in FCC) is subjected to a traction field of (100, 10, 10) MPa.



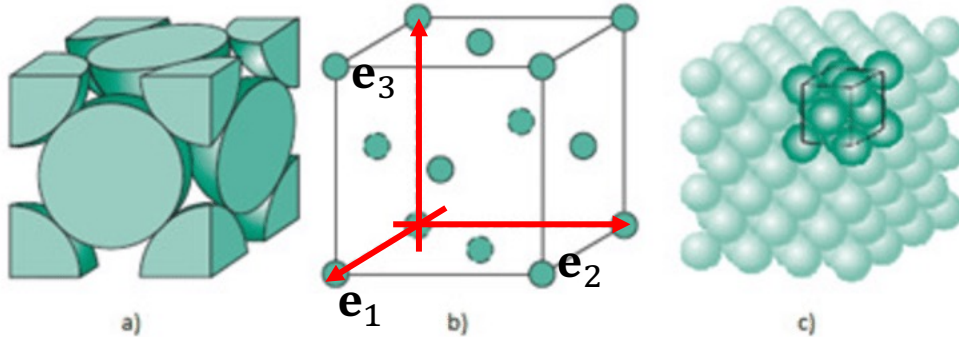
Calculate the traction vector along $(1\bar{1}1)$ plane normal

(Solution). Since, say, $\mathbf{n} = \frac{(1,1,1)}{|(1,1,1)|} = \frac{1}{\sqrt{3}}(1, \bar{1}, 1)$. The component of traction vector (given as $\mathbf{t} = (100, 10, 10)$) along \mathbf{n} is obtained by projecting \mathbf{t} along \mathbf{n} such that

$$\mathbf{t} \cdot \mathbf{n} = \frac{100}{\sqrt{3}}$$

Resolved shear stress

Suppose a single crystal Fe (in BCC) is subjected a stress of $\begin{bmatrix} 20 & 0 & 0 \\ 0 & 20 & 0 \\ 0 & 0 & 0 \end{bmatrix}$



Calculate the resolved shear stress acting on system $(011)[\bar{1}\bar{1}1]$

(Solution). Since, say, $\mathbf{n} = \frac{(0,1,1)}{|(0,1,1)|} = \frac{1}{\sqrt{2}}(0,1,1)$. The traction vector (denoted as \mathbf{t}) on plane \mathbf{n} is

$$\mathbf{t} = \frac{\sqrt{2}}{2} \begin{bmatrix} 20 & 0 & 0 \\ 0 & 20 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \\ 1 \end{bmatrix} = \frac{\sqrt{2}}{2} \begin{bmatrix} 20 \\ 0 \\ 0 \end{bmatrix}$$

In order to obtain \mathbf{t} component in the slip direction (denoted as \mathbf{b}), use dot-product, such

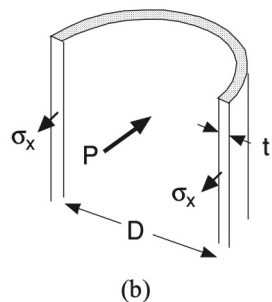
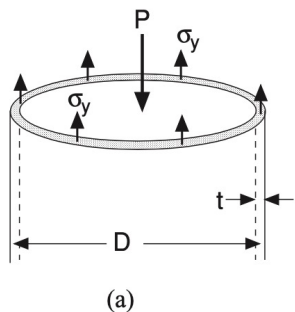
$$\text{that } \tau = \mathbf{t} \cdot \mathbf{b} = \frac{\sqrt{2}}{2} \frac{1}{\sqrt{3}} \begin{bmatrix} 20 \\ 0 \\ 0 \end{bmatrix} \cdot \begin{bmatrix} 1 \\ -1 \\ 1 \end{bmatrix} = \frac{\sqrt{2}}{2} \frac{1}{\sqrt{3}} \cdot 20 \cong 8.16$$

Generalized Schmid's law

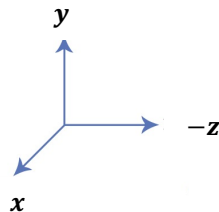
$$\tau = (\boldsymbol{\sigma} \cdot \boldsymbol{n}) \cdot \boldsymbol{b} = \boldsymbol{b} \cdot (\boldsymbol{\sigma} \cdot \boldsymbol{n}) = \boldsymbol{b} \cdot \boldsymbol{\sigma} \cdot \boldsymbol{n} = b_i \sigma_{ij} n_j$$

Force and moment balances

- Net force acting on any portion of a body should be zero.
- Externally applied forces (die, punch and so forth) should be balanced by internal forces – attractions and repulsions between atoms.



* 길이가 L 이고, 지름이 D 그리고 두께가 t인 파이프의 내부가 압력 P를 겪는다.



길이 방향이 normal인 면에 작용하는 응력 성분 σ_y 은, 이러한 force balance를 활용하여 구할 수 있다.

$$\sum_{\text{source}} F_y^{\text{each source}} = F_y^{\text{pipe pressure}} + F_y^{\text{internal stress}} = -\frac{P\pi D^2}{4} + \pi D t \sigma_y = 0$$

$$\rightarrow \pi D t \sigma_y = \frac{P\pi D^2}{4} \rightarrow \sigma_y = \frac{PD}{4t}$$

$$\sum_{\text{source}} F_x^{\text{each source}} = -PDL + 2\sigma_x tL = 0 \rightarrow \sigma_x = \frac{PD}{2t}$$

Force and moment balances

- $\mathbf{m} = \mathbf{F} \times \mathbf{x} \rightarrow \mathbf{m}_k = f_i x_j \mathbf{e}_i \times \mathbf{e}_j = f_i x_j \epsilon_{ijk} \mathbf{e}_k \quad \mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3 \parallel \hat{\mathbf{r}}, \hat{\boldsymbol{\theta}}, \hat{\mathbf{z}}$

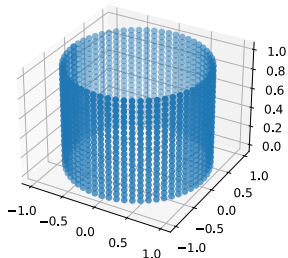
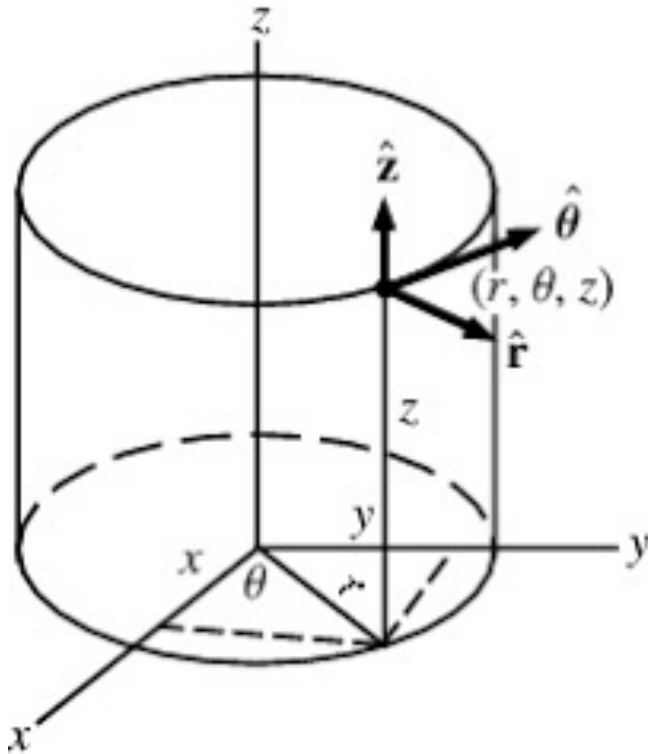
Shear stress

$$\sigma_{r\theta}$$

$$\mathbf{m}_3 = f_i x_j \epsilon_{ij3} \mathbf{e}_3 = (f_1 x_2 - f_2 x_1) \mathbf{e}_3$$

$$\mathbf{m}_3 = (t_1 A_r x_2 - t_2 A_\theta x_1) \mathbf{e}_3$$

$$= \left(\sigma_{1j} n_j^{(r)} A_r x_2 - \sigma_{2k} n_k^{(r)} A_\theta x_1 \right) \mathbf{e}_3$$



$$\sigma_{r\theta}$$

